

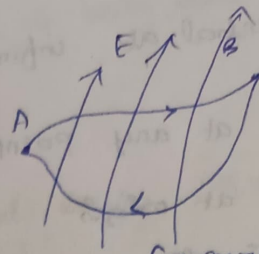
From the definition

$$V = -\int E \cdot dl$$

$$dV = -E \cdot dl$$

$$dV = -E_x dx - E_y dy - E_z dz$$

$$dV = \frac{\partial V}{\partial x} dx + \frac{\partial V}{\partial y} dy + \frac{\partial V}{\partial z} dz = -E \cdot dl$$



Conservative Nature  
of Electrostatic field

Comparing two expressions,

$$E_x = -\frac{\partial V}{\partial x}, \quad E_y = -\frac{\partial V}{\partial y}, \quad E_z = -\frac{\partial V}{\partial z}$$

$$\therefore \boxed{E = -\nabla V}$$

Electric field intensity is the gradient of  $V$ .

Negative sign shows that the direction of  $E$  is opposite to the direction in which  $V$  increases, i.e.  $E$  is directed from higher to lower levels of  $V$ .

Potential due to different charge distributions.

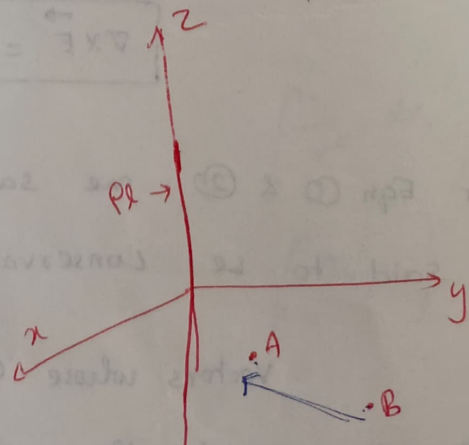
1) Uniformly charged line.

$$V_{ba} = -\int_b^a E \cdot dl$$

$$= -\int_b^a E \cdot dr$$

$$= -\frac{\rho l}{2\pi\epsilon_0} \int_b^a \frac{1}{r} dr$$

$$= -\frac{\rho l}{2\pi\epsilon_0} [\ln(r)]_b^a$$

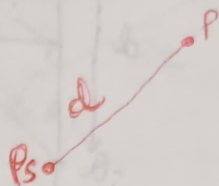


$$= -\frac{\rho l}{2\pi\epsilon_0} (\ln a - \ln b)$$

$$= \frac{\rho l}{2\pi\epsilon_0} (\ln b - \ln a)$$

$$V_{ba} = \frac{\rho l}{2\pi\epsilon_0} \ln(b/a) \quad \checkmark$$

ii) Charged disc



$$V = - \int E \cdot dl$$

$$= - \int_d^0 E \cdot dx$$

$$= - \int_d^0 \frac{\rho_s}{2\epsilon} (1 - \cos\alpha) dx$$

$$= \frac{\rho_s}{2\epsilon} (1 - \cos\alpha) \int_d^0 dx$$

$$= \frac{\rho_s}{2\epsilon} (1 - \cos\alpha) (0 - d)$$

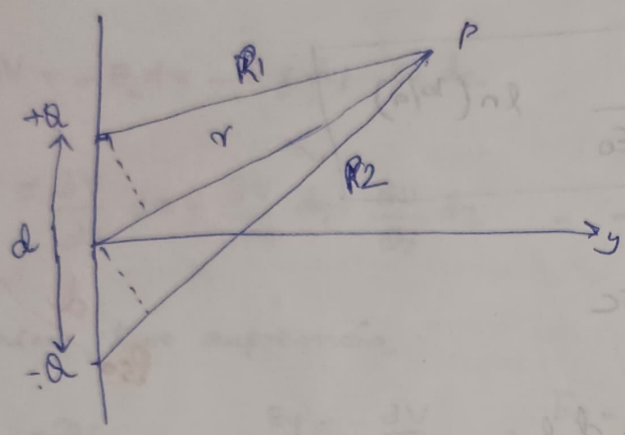
$$= \frac{\rho_s}{2\epsilon} (1 - \cos\alpha) d$$

$$= \frac{\rho_s}{2\epsilon} \left(1 - \frac{h}{\sqrt{h^2 + R^2}}\right) \sqrt{h^2 + R^2}$$

$$V = \frac{\rho_s}{2\epsilon} (\sqrt{h^2 + R^2} - h)$$

Potential due to Electric dipole

Equal and opposite <sup>point</sup> charges separated by a small distance is known as electric dipole.



Potential at P due to +Q is

$$V_1 = \frac{Q}{4\pi\epsilon_0 R_1}$$

Potential at P due to -Q is

$$V_2 = -\frac{Q}{4\pi\epsilon_0 R_2}$$

Total Potential at P due to the dipole is

$$V = V_1 + V_2$$

$$= \frac{Q}{4\pi\epsilon_0 R_1} - \frac{Q}{4\pi\epsilon_0 R_2}$$

$$= \frac{Q}{4\pi\epsilon_0} \left[ \frac{1}{R_1} - \frac{1}{R_2} \right]$$

$$V = \frac{Q}{4\pi\epsilon_0} \left[ \frac{R_2 - R_1}{R_1 R_2} \right]$$

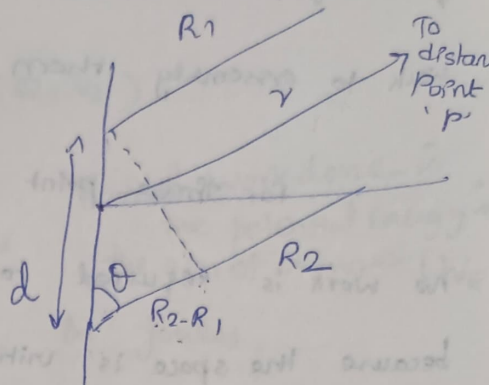
for a distant point

$$R_1 \approx R_2 = r$$

$$R_1 R_2 = r^2$$

(for the numerator) if  $R_1$  &  $R_2$  are assumed to be parallel,

then  $R_2 - R_1 = d \cos \theta$



$$V = \frac{Q}{4\pi\epsilon_0} \left( \frac{d \cos \theta}{r^2} \right)$$

$$V = \frac{Q d \cos \theta}{4\pi\epsilon_0 r^2}$$

Dipole moment  $m = Qd$

Potential:

$$V = \frac{m \cos \theta}{4\pi\epsilon_0 r^2}$$

Electric field intensity

$$E = -\nabla V$$

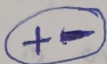
Assume spherical coordinate system

$$E = - \left[ \frac{\partial V}{\partial r} \hat{a}_r + \frac{1}{r} \frac{\partial V}{\partial \theta} \hat{a}_\theta + \frac{1}{r \sin \theta} \frac{\partial V}{\partial \phi} \hat{a}_\phi \right]$$

$$= - \left[ \frac{m \cos \theta}{4\pi\epsilon_0} \left( \frac{-2}{r^3} \right) \hat{a}_r + \frac{1}{r} \frac{m}{4\pi\epsilon_0 r^2} (-\sin \theta) \hat{a}_\theta + \frac{1}{r \sin \theta} (0) \hat{a}_\phi \right]$$

$$= \frac{m \cos \theta}{2\pi\epsilon_0 r^3} \hat{a}_r + \frac{m \sin \theta}{4\pi\epsilon_0 r^3} \hat{a}_\theta$$

$$E = \frac{m}{4\pi\epsilon_0 r^3} \left[ 2 \cos \theta \hat{a}_r + \sin \theta \hat{a}_\theta \right]$$



$$\frac{d^2}{dx^2} x^{n-1} = \frac{d}{dx} (n-1)x^{n-2} = (n-1)(n-2)x^{n-3}$$

$$\left[ \begin{matrix} -2r^{-2} \\ -2r^{-3} \\ -2r^{-3} \end{matrix} \right]$$