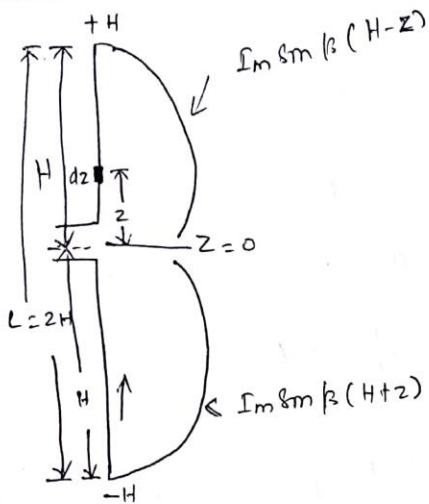
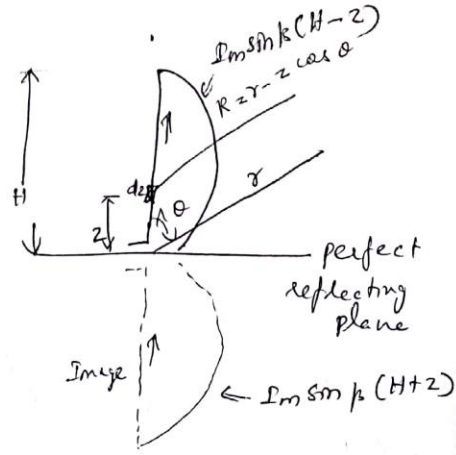




Radiation from half wave dipole and quarter-wave monopole



(a) centre-fed dipole with assumed sinusoidal current distribution



(b) corresponding monopole

Figure (a) shows a centre-fed dipole with a sinusoidal current distribution & (b) shows corresponding monopole.

"A dipole antenna" is a straight radiator, usually fed in the center and producing a maximum of radiation in the plane normal to the axis. The length specified is the overall length.

Radiation from a quarter-wave monopole or halfwave dipole

It will be assumed that the current is sinusoidally distributed as shown in figure.

Then,

$$I = I_m \sin \beta (H - z) \quad , \quad z \geq 0$$

$$I = I_m \sin \beta (H + z) \quad , \quad z < 0$$

$I_m \rightarrow$ The value of current at the current loop (or) current maximum.



The expression for the vector potential at point P due to the current element $I dz$ is

$$dA_z = \frac{\mu I e^{-j\beta R}}{4\pi R} dz \rightarrow \textcircled{1}$$

$R \rightarrow$ distance from the current element to point P. The total vector potential at P due to all the current elements will be,

$$A_z = \frac{\mu}{4\pi} \int_{-H}^0 \frac{I_m \sin \beta (H+z) e^{-j\beta R}}{R} dz + \frac{\mu}{4\pi} \int_0^H \frac{I_m \sin \beta (H-z) e^{-j\beta R}}{R} dz \rightarrow \textcircled{2}$$

~~$R \approx r$~~ $\boxed{R \approx r}$ \rightarrow (in denominator)

$R \approx r$ \rightarrow R in the phase factor, so the difference between R and r is important.

$$R = r - z \cos \theta.$$

Then the expression for A_z becomes,

$$A_z = \frac{\mu I_m e^{-j\beta r}}{4\pi r} \left[\int_{-H}^0 \sin \beta (H+z) e^{j\beta z \cos \theta} dz + \int_0^H \sin \beta (H-z) e^{j\beta z \cos \theta} dz \right] \rightarrow \textcircled{3}$$

[if $H = \lambda/4$.

$$\begin{aligned} \sin \beta (H+z) &= \sin \beta (H-z) \\ &= \sin(\pi/2 + \beta z) = \sin(\pi/2 - \beta z) \\ &= \cos \beta z. \end{aligned}$$

$$\frac{\beta H}{\pi} \times \frac{\pi}{\lambda} = \frac{\pi}{4}$$

Then eq³ becomes,

$$A_z = \frac{\mu I_m e^{-j\beta r}}{4\pi r} \left[\int_{-H}^0 \cos \beta z e^{j\beta z \cos \theta} dz + \int_0^H \cos \beta z e^{j\beta z \cos \theta} dz \right]$$



$$\begin{aligned}
 &= \frac{\mu I_m e^{-j\beta r}}{4\pi r} \left[\int_0^H \cos \beta z e^{-j\beta z \cos \theta} dz + \int_0^H \cos \beta z e^{j\beta z \cos \theta} dz \right] \\
 &= \frac{\mu I_m e^{-j\beta r}}{4\pi r} \int_0^H \cos \beta z (e^{j\beta z \cos \theta} + e^{-j\beta z \cos \theta}) dz \\
 &= \frac{\mu I_m e^{-j\beta r}}{4\pi r} \int_0^H 2 \cos \beta z \cos \beta z \cos \theta dz \\
 &= \frac{\mu I_m e^{-j\beta r}}{4\pi r} \int_0^H [\cos(\beta z + \beta z \cos \theta) + \cos(\beta z - \beta z \cos \theta)] dz \\
 &= \frac{\mu I_m e^{-j\beta r}}{4\pi r} \int_0^H [\cos \beta z (1 + \cos \theta) + \cos \beta z (1 - \cos \theta)] dz \\
 &= \frac{\mu I_m e^{-j\beta r}}{4\pi r} \left[\frac{\sin \beta z (1 + \cos \theta)}{\beta (1 + \cos \theta)} + \frac{\sin \beta z (1 - \cos \theta)}{\beta (1 - \cos \theta)} \right]_0^H \\
 &= \frac{\mu I_m e^{-j\beta r}}{4\pi r \beta} \left[\frac{(1 - \cos \theta) \cos \frac{\pi}{2} \cos \theta + (1 + \cos \theta) \cos \frac{\pi}{2} \cos \theta}{\sin^2 \theta} \right] \\
 &= \frac{\mu I_m e^{-j\beta r}}{4\pi r \beta} \left[\frac{\cos \frac{\pi}{2} \cos \theta (1 - \cos \theta + 1 + \cos \theta)}{\sin^2 \theta} \right] \\
 &= \frac{\mu I_m e^{-j\beta r}}{2\pi r \beta} \times \frac{2 \cos \frac{\pi}{2} \cos \theta}{\sin^2 \theta} \\
 \boxed{A_z = \frac{\mu I_m e^{-j\beta r}}{2\pi \beta r} \left[\frac{\cos \frac{\pi}{2} \cos \theta}{\sin^2 \theta} \right]} &\rightarrow (4)
 \end{aligned}$$

when current is entirely in z direction,

$$\mu H_\phi = \frac{1}{r} \frac{\partial}{\partial r} (r A_z \sin \theta)$$

~~$$\mu H_\phi = \frac{\partial A_z \sin \theta}{\partial r} \rightarrow (5)$$~~

$$H_\phi = \frac{1}{\mu_0} \left[-\frac{\partial}{\partial r} \times \frac{\mu I_m e^{-j\beta r}}{2\pi \beta r} \frac{\cos \frac{\pi}{2} \cos \theta}{\sin^2 \theta} \right]$$



$$= \frac{1}{4\pi r} \left[\frac{j I_m e^{-jkr}}{2\pi r} \frac{\cos(\pi/2 \cos \theta)}{\sin^2 \theta} \right] \times (r d\Omega) / r$$

$$H_\phi = \frac{j I_m}{2\pi r} e^{-jkr} \left[\frac{\cos(\pi/2 \cos \theta)}{\sin^2 \theta} \right] \rightarrow (6)$$

The electric field strength at a distant point will be,

$$E_\theta = \eta H_\phi$$

$$= 120 \pi \times \frac{j I_m}{2\pi r} e^{-jkr} \left[\frac{\cos(\pi/2 \cos \theta)}{\sin^2 \theta} \right]$$

$$E_\theta = \frac{j 60 I_m}{r} e^{-jkr} \left[\frac{\cos(\pi/2 \cos \theta)}{\sin^2 \theta} \right] \rightarrow (7)$$

The magnitude of the electric field strength for the radiation field of half-wave dipole (or) quarter-wave monopole is

$$|E_\theta| = \frac{60 I_m}{r} \left[\frac{\cos(\pi/2 \cos \theta)}{\sin \theta} \right] \text{ V/m.} \rightarrow (8)$$

Max. value
Poynting vector

$$P_{max} = |E_\theta| \times |H_\phi|$$

$$= \frac{60 I_m}{r} \frac{\cos(\pi/2 \cos \theta)}{\sin \theta} \times \frac{I_m}{2\pi r} \frac{\cos(\pi/2 \cos \theta)}{\sin^2 \theta}$$

$$P_{av} = P_{max}$$

$$= \frac{15}{\pi} \frac{I_m^2}{r^2} \frac{\cos(\pi/2 \cos \theta)}{\sin \theta} \times \frac{I_m}{2\pi r} \frac{\cos(\pi/2 \cos \theta)}{\sin \theta}$$

$$P_{av} = \frac{15 I_m^2}{\pi r^2} \left\{ \frac{\cos(\pi/2 \cos \theta)}{\sin \theta} \right\}^2$$

In terms of rms current

$$I_{rms} = \frac{I_m}{\sqrt{2}}$$

$$I_m = \sqrt{2} I_{rms}$$

