

17/05/23

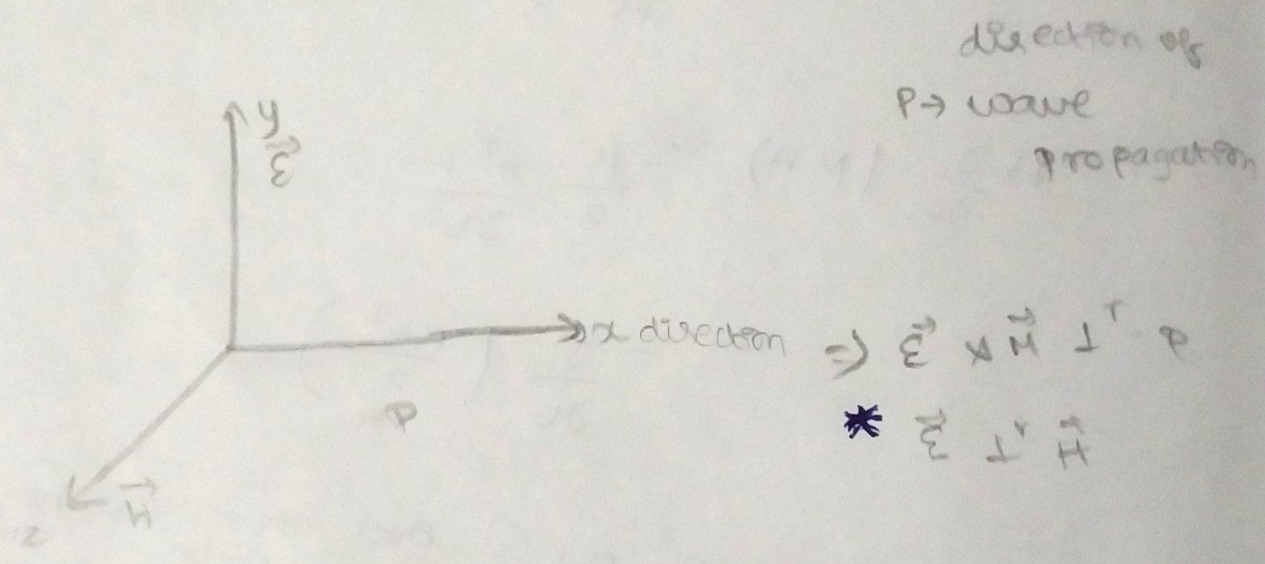
1st period

Unit - \vec{v}

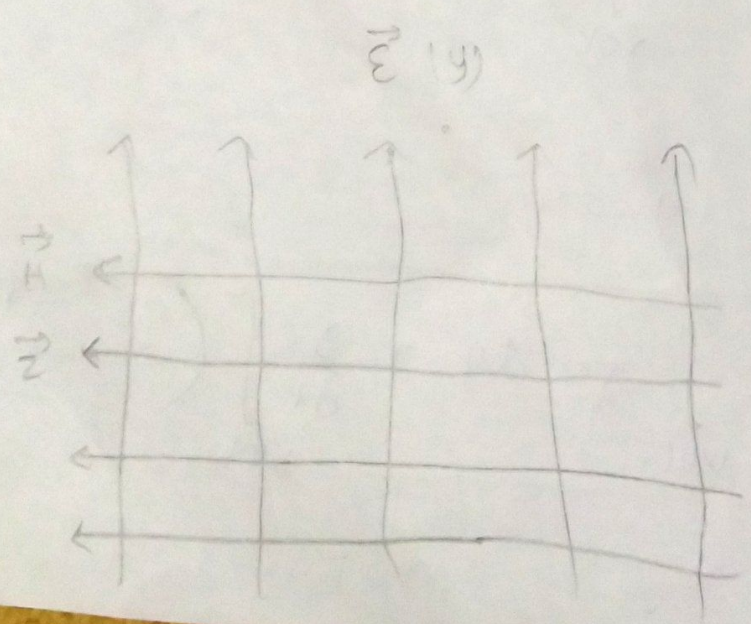
Electromagnetic Waves

like spherical wave but r is

Derivation of wave equation - Uniform Plane waves



* Also called as Transverse Electromagnetic wave.



orientation of polarization of elec. field at any instant of time.

* vertically polarized.

Maxwell's eqn from ampere's law.

$$\nabla \times H = J + \frac{\partial D}{\partial t} \quad \text{--- (1)}$$

for free space ($J = \sigma E$),
 $\sigma = 0, J = 0$

$$\nabla \times H = \frac{\partial D}{\partial t}$$

$$[\because D = \epsilon_0 E]$$

$$\nabla \times H = \epsilon_0 \frac{\partial E}{\partial t} \quad \text{--- (2)}$$

Expanding (2)

$$\begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ H_x & H_y & H_z \end{vmatrix} = \frac{\partial}{\partial t} (\hat{x} \mathcal{D}_x + \hat{y} \mathcal{D}_y + \hat{z} \mathcal{D}_z)$$

$$\hat{x} \left[\frac{\partial H_z}{\partial y} - \frac{\partial H_y}{\partial z} \right] - \hat{y} \left[\frac{\partial H_z}{\partial x} - \frac{\partial H_x}{\partial z} \right] + \hat{z} \left[\frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} \right]$$

$$= \frac{\partial}{\partial t} (\hat{x} \mathcal{D}_x + \hat{y} \mathcal{D}_y + \hat{z} \mathcal{D}_z)$$

\therefore the wave is propagating in x direction, the

only component contribute is,

$$-\hat{y} \left(\frac{\partial H_z}{\partial x} \right) = \hat{y} \frac{\partial \mathcal{D}_y}{\partial t} \quad \text{--- (3)}$$

$$\text{flux } D = \epsilon E$$

$$\frac{\partial H_z}{\partial x} = -\epsilon \frac{\partial E_y}{\partial t} \quad \text{--- (4)}$$

Eqn (4) relates space derivative of H_z to time derivative of E_y .

Maxwell's second eqn from Faraday's law.

$$\nabla \times E = -\frac{\partial B}{\partial t}$$

$$\begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ E_x & E_y & E_z \end{vmatrix} = -\frac{\partial}{\partial t} (\hat{x} B_x + \hat{y} B_y + \hat{z} B_z)$$

$$\hat{x} \left(\frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} \right) - \hat{y} \left(\frac{\partial E_z}{\partial x} - \frac{\partial E_x}{\partial z} \right) + \hat{z} \left(\frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} \right)$$

$$= -\frac{\partial}{\partial t} (\hat{x} B_x + \hat{y} B_y + \hat{z} B_z)$$

Since the wave propagating in x direction, the only component contribute is,

$$\hat{z} \left[\frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} \right] = -\hat{z} \frac{\partial B_z}{\partial t} \quad \text{--- (7)}$$

$$\hat{z} \left[\frac{\partial E_y}{\partial x} \right] = -\hat{z} \frac{\partial B_z}{\partial t}$$

$$\frac{\partial E_y}{\partial x} = -\mu \frac{\partial H_z}{\partial t} \quad // \quad \textcircled{8}$$

eqn ⑧ relates space derivative of E_y to time derivative of H_z .

18/07/23
* period

Differentiating eq ⑧ wrt time.

$$\frac{\partial}{\partial t} \left(\frac{\partial H_z}{\partial x} \right) = -\epsilon_0 \frac{\partial}{\partial t} \left(\frac{\partial E_y}{\partial x} \right)$$

\downarrow time derivative \downarrow space derivative

$$\frac{\partial}{\partial t} \left(\frac{\partial H_z}{\partial x} \right) = -\epsilon_0 \frac{\partial^2 E_y}{\partial t^2} \quad // \quad \textcircled{9}$$

eqn ⑧ wrt distance 'x'

$$\frac{\partial}{\partial x} \left(\frac{\partial E_y}{\partial x} \right) = -\mu \frac{\partial}{\partial x} \left(\frac{\partial H_z}{\partial t} \right)$$

$$\frac{\partial^2 E_y}{\partial x^2} = -\mu \frac{\partial}{\partial x} \left(\frac{\partial H_z}{\partial t} \right) \quad // \quad \textcircled{10}$$

÷ eq ⑩ by $-\mu$.

$$-\frac{1}{\mu} \left(\frac{\partial^2 E_y}{\partial x^2} \right) = \frac{\partial}{\partial x} \left(\frac{\partial H_z}{\partial t} \right) \quad // \quad \textcircled{11}$$

Compare ⑨ & ⑪

LHS of ⑨ = RHS of ⑪

$$-\epsilon_0 \frac{\partial^2 E_y}{\partial t^2} = -\frac{1}{\mu} \left(\frac{\partial^2 E_y}{\partial x^2} \right)$$

$$\frac{\partial^2 \epsilon_y}{\partial t^2} = \frac{1}{\mu \epsilon} \left[\frac{\partial^2 \epsilon_y}{\partial x^2} \right] \quad \text{--- (12)}$$

Laplace eq. respect to ϵ_y

Diff eq (12) w.r.t distance

$$\frac{\partial}{\partial x} \left(\frac{\partial H_z}{\partial x} \right) = -\epsilon \frac{\partial}{\partial x} \left(\frac{\partial \epsilon_y}{\partial t} \right)$$

$$\frac{\partial^2 H_z}{\partial x^2} = -\epsilon \frac{\partial}{\partial x} \left(\frac{\partial \epsilon_y}{\partial t} \right) \quad \text{--- (13)}$$

Diff eq (13) w.r.t time

$$\frac{\partial}{\partial t} \left(\frac{\partial \epsilon_y}{\partial x} \right) = -\mu \frac{\partial}{\partial t} \left(\frac{\partial H_z}{\partial x} \right)$$

$$\frac{\partial}{\partial t} \left(\frac{\partial \epsilon_y}{\partial x} \right) = -\mu \frac{\partial^2 H_z}{\partial t^2} \quad \text{--- (14)}$$

Divide eq (14) by $-\epsilon$

$$-\frac{1}{\epsilon} \left(\frac{\partial}{\partial t} \right) \left(\frac{\partial \epsilon_y}{\partial x} \right) = \frac{\epsilon \mu}{1} \frac{\partial^2 H_z}{\partial t^2} \quad \text{--- (15)}$$

Compare (13) & (15)

$$\frac{\partial^2 H_z}{\partial x^2} = \epsilon \mu \left(\frac{\partial^2 H_z}{\partial t^2} \right) \quad \text{--- (16)}$$

$$\frac{\partial^2 H_z}{\partial t^2} = \frac{1}{\mu \epsilon} \left[\frac{\partial^2 H_z}{\partial x^2} \right] \quad \text{--- (16)}$$

eq (12) & (16) are also called as $\nabla \cdot \mathbf{A} = \mu_0 \rho$ & $\nabla \times \mathbf{A} = \mu_0 \mathbf{j}$ equations

$$\frac{1}{\mu \epsilon} = \gamma$$

velocity (v) depends on the parameters μ & ϵ

Q. In a material for which conductivity $\sigma = 4.5 \text{ v/m}$
in a material for which conductor $\epsilon_r = 1$.
and the electric field $\mathbf{E} = 300 \text{ qsn } 10^9 \hat{t} \text{ v/m}$. Determine
the conduction & displacement current densities
and also find the frequency at which they have
equal magnitudes.

soln:

Given, $\sigma = 4.5 \text{ v/m}$

$$\epsilon_r = 1$$

$$\mathbf{E} = 300 \text{ qsn } 10^9 \hat{t} \text{ v/m}$$

$$\mathbf{J}_c = \sigma \times \mathbf{E}$$

$$= 4.5 \times 300 \text{ qsn } 10^9 \hat{t}$$

$$\mathbf{J}_c = 1350 \text{ qsn } 10^9 \hat{t} \text{ A/m}^2 //$$

$$\frac{\partial \mathbf{D}}{\partial t} = \epsilon \frac{\partial \mathbf{E}}{\partial t}$$

$$= \epsilon \frac{\partial (300 \text{ qsn } 10^9 \hat{t})}{\partial t}$$

$$= \epsilon_0 \epsilon_r \frac{\partial (300 \text{ qsn } 10^9 \hat{t})}{\partial t}$$

$$\frac{\partial D}{\partial t} = 8.854 \times 10^{-12} \frac{\partial}{\partial t} (300 \sin 10^9 t)$$

$$= 8.854 \times 10^{-12} \cos 10^9 t (10^9)(300)$$

$$= 8.854 \times 10^{-3} \cos 10^9 t (300)$$

$$= 2.656 \cos 10^9 t \text{ // } = 2.652 \times 10^{-10} \text{ A/m}^2$$

28/05/23

3rd period.

wave Equations And

Their Solutions

Maxwell's second eqn in pointed form

$$\nabla \times \mathcal{E} = - \frac{\partial \mathcal{B}}{\partial t}$$

$$\nabla \times \mathcal{E} = - \mu \frac{\partial \mathcal{H}}{\partial t} \quad \text{--- (1)}$$

Take curl on both sides of eq (1)

$$\nabla \times \nabla \times \mathcal{E} = - \mu \left[\nabla \times \frac{\partial \mathcal{H}}{\partial t} \right] \quad \text{--- (2)}$$

Maxwell's first eqn in point form

$$\nabla \times \mathcal{H} = \mathcal{J} + \frac{\partial \mathcal{D}}{\partial t}$$

$$\nabla \times \mathcal{H} = \sigma \mathcal{E} + \frac{\partial \mathcal{D}}{\partial t} \quad \text{--- (3)}$$

Differentiating eq (7)

$$\nabla \times \frac{\partial \mathbf{H}}{\partial t} = \sigma \frac{\partial \mathbf{E}}{\partial t} + \frac{\partial^2 \mathbf{D}}{\partial t^2} \quad \text{--- (6)}$$

sub eq (6) in (5)

$$\nabla \times \nabla \times \mathbf{E} = -\mu \left[\sigma \frac{\partial \mathbf{E}}{\partial t} + \frac{\partial^2 \mathbf{D}}{\partial t^2} \right] \quad \text{--- (5)}$$

Using an identity,

$$\nabla \times \nabla \times \mathbf{E} = \nabla (\nabla \cdot \mathbf{E}) - \nabla^2 \mathbf{E}$$

$$\nabla \times \mathbf{E} = \nabla \cdot \frac{\mathbf{D}}{\epsilon}$$

$$= \frac{1}{\epsilon} [\nabla \cdot \mathbf{D}] \quad \text{--- (6)}$$

Material is taken as a conductor, ($\rho = 0$) \rightarrow within a conductor

$$\nabla \cdot \mathbf{D} = \rho$$

$$\nabla \cdot \mathbf{D} = 0 //$$

$$\therefore \nabla \cdot \mathbf{E} = 0 //$$

Using identity,

$$\nabla \times \nabla \times \mathbf{E} = -\nabla^2 \mathbf{E} \quad \text{--- (7)}$$

Compare eq (5) & (6)

$$\nabla \times \nabla \times \mathbf{E} = -\nabla^2 \mathbf{E}$$

$$-\nabla^2 \mathbf{E} = - \left[\mu \sigma \frac{\partial \mathbf{E}}{\partial t} + \mu \frac{\partial^2 \mathbf{D}}{\partial t^2} \right]$$

$$\nabla^2 E = \mu \sigma \frac{\partial E}{\partial t} - \mu \epsilon \frac{\partial^2 E}{\partial t^2} = 0 \quad \text{--- (8)}$$

Eqn (8) is the wave equation for electric field.

Derive wave eqn in terms of mag. field.

Maxwell's first eqn in point form.

$$\nabla \times H = J + \frac{\partial D}{\partial t}$$

$$\nabla \times H = \sigma E + \epsilon \frac{\partial E}{\partial t} \quad \text{--- (9)}$$

Diff eq (9)

Take curl on both sides,

$$\nabla \times \nabla \times H = \sigma \nabla \times E + \epsilon \nabla \times \frac{\partial E}{\partial t}$$

Maxwell's 2nd eqn in point form.

$$\nabla \times E = -\frac{\partial B}{\partial t}$$

$$\nabla \times E = -\mu \frac{\partial H}{\partial t} \quad \text{--- (10)}$$

Diff (10) w.r.t time

$$\nabla \times \frac{\partial E}{\partial t} = -\mu \frac{\partial^2 H}{\partial t^2} \quad \text{--- (11)}$$

$$\nabla \times \frac{\partial \mathbf{A}}{\partial t} = \nabla \times \mathbf{A} \quad \text{--- (1)}$$

$$\nabla \times \nabla \times \mathbf{H} = \sigma \nabla \cdot \mathbf{E} + \mathbf{E} \left[-\mu \frac{\partial^2 \mathbf{H}}{\partial t^2} \right] \quad \text{--- (2)}$$

Using an vector identity,

$$\nabla \times \nabla \times \mathbf{H} = \nabla (\nabla \cdot \mathbf{H}) - \nabla^2 \mathbf{H}$$

$$\nabla \times \nabla \times \mathbf{H} = 0$$

$$\nabla \times \mathbf{H} = \nabla \cdot \frac{\mathbf{B}}{\mu_0}$$

$$\nabla \times \mathbf{H} = (\sigma + j\omega\epsilon) \mathbf{E}$$

$$\nabla \times \mathbf{H} = \frac{1}{\mu} [\nabla \times \mathbf{B}] \quad \text{--- (3)}$$

within a conductor,

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \cdot \mu \mathbf{H} = 0$$

$$\nabla \cdot \mathbf{H} = 0$$

Sub $\nabla \cdot \mathbf{H}$

$$\nabla \times \nabla \times \mathbf{H} = \nabla (0) - \nabla^2 \mathbf{H}$$

$$\nabla \times \nabla \times \mathbf{H} = -\nabla^2 \mathbf{H}$$

$$-\nabla^2 H = \sigma \left(-\mu \frac{\partial H}{\partial t} \right) + \epsilon \left[-\mu \frac{\partial^2 H}{\partial t^2} \right]$$

$$+\nabla^2 H = \sigma \mu \frac{\partial H}{\partial t} + \epsilon \mu \left[\frac{\partial^2 H}{\partial t^2} \right]$$

$$\nabla^2 H - \sigma \mu \left[\frac{\partial H}{\partial t} \right] + \epsilon \mu \left[\frac{\partial^2 H}{\partial t^2} \right] = 0$$

Wave equation in terms of mag. field.

25/5/23

3rd period.

$$\gamma = \alpha + j\beta$$

(i.e) Propagation constant = ^{Attenuation} ~~attenuation~~ constant + Phase constant

Case i)

$$\gamma \rightarrow \text{real } [\gamma = \alpha], \beta = 0.$$

↓
No wave propagation

Case ii)

$$\gamma \rightarrow \text{imaginary } (\gamma = j\beta), \alpha = 0.$$

* In high frequencies used.

Wave equation in Phasor form:

Maxwell equation in Phasor form:

Maxwell's first equation,

$$\nabla \times H = J + \frac{\partial D}{\partial t}$$

$$\left[\because \frac{\partial D}{\partial t} = j\omega D \right]$$

$$\nabla \times H = \sigma E + j\omega \epsilon E$$

$$\boxed{\nabla \times H = (\sigma + j\omega \epsilon) E}$$

$$\nabla \times E = -\frac{\partial B}{\partial t}$$

$$\boxed{\nabla \times E = -j\omega \mu H}$$

Wave eqn in E

$$\nabla^2 E - \mu \sigma \frac{\partial E}{\partial t} - \mu \epsilon \frac{\partial^2 E}{\partial t^2} = 0$$

$$\nabla^2 E - \mu \sigma (j\omega) - \mu \epsilon (j\omega)^2 E = 0$$

$$\nabla^2 E - \frac{j\omega \mu (\sigma + j\omega \epsilon E)}{1} = 0$$

$$\nabla^2 E - \gamma^2 E = 0$$

$$\Rightarrow \nabla^2 E = \gamma^2 E$$

1/4

$$\nabla^2 H - \gamma^2 H = 0$$

$$\Rightarrow \nabla^2 H = \gamma^2 H$$

$$\gamma^2 = j\omega\mu [\sigma + j\omega\epsilon E]$$

$$\gamma = \sqrt{j\omega\mu (\sigma + j\omega\epsilon E)}$$

$$\gamma = \alpha + j\beta //$$

$$\alpha = \sqrt{\frac{\mu\epsilon}{2} \left[\left(\sqrt{\left(1 + \frac{\sigma^2}{\omega^2\epsilon^2}\right)^2} - 1 \right) \right]}$$

neper / metre //

$$\beta = \sqrt{\frac{\mu\epsilon}{2} \left[\left(\sqrt{\left(1 + \frac{\sigma^2}{\omega^2\epsilon^2}\right)^2} + 1 \right) \right]}$$

radus / metre //

Impedance impedance:

$$\eta = \sqrt{\frac{j\omega\mu}{\sigma + j\omega\epsilon}} \quad \text{ohms} //$$

Ratio b/w electric field and magnetic field.

Free space $\sigma = 0$

$$\eta_0 = \sqrt{\frac{j\omega\mu}{j\omega\epsilon}} = \sqrt{\frac{\mu}{\epsilon}}$$

$$\eta_0 = \sqrt{\frac{\mu_0}{\epsilon_0}} \quad \text{ohms} //$$

$$\eta_0 = 120\pi \text{ (ohms)} \approx 377 \Omega$$

wavelength (λ)

$$\lambda = \frac{v}{f} = \frac{2\pi}{\beta}$$

Free space $v = c$

$$\lambda = \frac{c}{f}$$

① A 6580 MHz uniform plane wave is propagating in a material medium of $\epsilon_r = 2.25$. If the amplitude of electric field intensity of lossless medium is 500 Volts/m. Calculate

- i) Phase constant β
- ii) Propagation constant
- iii) Phase velocity
- iv) Wavelength
- v) Intrinsic impedance
- vi) Amplitude of mag. field intensity.

Soln:-

$$\begin{aligned} \text{i) } \beta &= \omega \sqrt{\mu \epsilon} \\ &= 2\pi f \sqrt{\mu_0 \mu_r \epsilon_0 \epsilon_r} \\ &= 2 \times 3.14 \times 6580 \times 10^6 \sqrt{4\pi \times 10^{-7} \times 1 \times 8.854 \times 10^{-12} \times 2.25} \\ &= 206.859 \text{ rad/m} \end{aligned}$$

$$\begin{aligned} \text{ii) } \gamma &= \alpha + j\beta \\ &= 0 + j(206.859) \end{aligned}$$

$$ii) v_p = \frac{\omega}{\beta}$$

$$= \frac{2 \times 3.14 \times 6580 \times 10^6}{206.859}$$

$$= 200 \times 10^6 \text{ m/s} //$$

$$= 200 \times 10^6 \text{ m/s} //$$

iv)

$$\lambda = \frac{2\pi}{\beta}$$

$$= \frac{2 \times 3.14}{206.859}$$

$$= 0.030$$

v)

$$\eta = \sqrt{\frac{j\omega\mu}{\sigma + j\omega\epsilon}}$$

$$\eta = \sqrt{\frac{\mu}{\epsilon_2}}$$

$$= \frac{4\pi \times 10^{-7} \times 1}{\sqrt{8.854 \times 10^{-12} \times 2.25}}$$

=

vi)

$$\epsilon = 500 \text{ v/m.}$$

$$\eta = \frac{\epsilon}{h}$$

$$= \frac{E}{h}$$

$$= \frac{500}{251.156}$$

$$= 1.99 \text{ Ampere/meter}$$

Plane wave ~~in~~ conductor

$$\frac{\sigma}{\omega \epsilon} \gg 1$$

$$\sigma \gg \omega \epsilon$$

$$\alpha = \omega \sqrt{\frac{\mu \epsilon}{2} \sqrt{1 + \left(\frac{\sigma}{\omega \epsilon}\right)^2} - 1}$$

$$= \omega \sqrt{\frac{\mu \epsilon}{2} \left[\frac{\sigma}{\omega \epsilon}\right]^2}$$

$$= \omega \sqrt{\frac{\mu \epsilon}{2} \times \frac{\sigma}{\omega \epsilon}}$$

$$= \sqrt{\frac{\omega^2 \mu \sigma}{2 \omega}}$$

$$= \sqrt{\frac{\omega \mu \sigma}{2}}$$

$$= \sqrt{\frac{2 \pi f \mu \sigma}{2}}$$

$$\alpha = \sqrt{\pi f \mu \sigma} \text{ //}$$

$$\therefore \alpha = \sqrt{\pi f \mu \sigma} \text{ nepers/m //}$$

$$\therefore \beta = \sqrt{\pi f \mu \sigma} \text{ radians/m //}$$

$$\left[\begin{array}{l} \text{proof:} \\ \beta = \sqrt{\epsilon \mu} \omega \end{array} \right]$$

Final result
be same by
elementary //

$$\eta = \sqrt{\frac{j\omega\mu}{\sigma + j\omega\epsilon}}$$

$$\eta = \sqrt{\frac{j\omega\mu}{\sigma}}$$

$$\eta = \sqrt{\frac{\omega\mu}{\sigma}} \angle 45^\circ \Omega //$$

($\frac{E}{H}$)

$\rightarrow E$ leads H by 45°

$\sigma \gg \omega\epsilon$, $\sigma + j\omega\epsilon$ will be neglected.
Because conductor value will be high.

Plane wave is conductor

$$v = \frac{\omega}{\beta} = \frac{2\pi f}{\sqrt{\pi f \mu \sigma}}$$

$$= \frac{\sqrt{2\pi} \times \sqrt{2\pi} \times \sqrt{f} \times \sqrt{f}}{\sqrt{\pi f \mu \sigma}}$$

$$= \frac{2\pi\sqrt{f}}{\sqrt{\mu\sigma}}$$

$$v = \sqrt{\frac{2\omega}{\mu\sigma}} //$$

Skin depth (or) depth of penetration (δ)

The E or H of the wave travels in a conducting medium its amplitude is attenuated by $e^{-\alpha z}$. The distance z through which the wave amplitude is attenuated by the

factor $e^{-\alpha z}$.

The distance through which the wave amplitude decreases by the factor e^{-1} (about 37% of the original value) is called skin depth or penetration depth of the medium.

$$\delta = \frac{1}{\alpha}$$

Hint: Skin depth is the measure of depth to which an electromagnetic wave can penetrate the medium.

① A uniform plane wave in a medium having conductivity $\sigma = 10^{-3} \text{ S/m}$, $\epsilon = 80 \epsilon_0$, $\mu = \mu_0$ is having a frequency of 10 kHz.

- i) Verify the medium is a good conductor or not.
- ii) Calculate α , β , δ , η , λ , v .

Soln:-

$$\frac{\sigma}{\omega \epsilon} \gg 1$$

$$\frac{\sigma}{\omega \epsilon} = \frac{10^{-3}}{2\pi \times 10^3 \times (80 \epsilon_0)} = \frac{10^{-3}}{2\pi \times 10^3 \times 8.854 \times 10^{-12}}$$

$$\alpha = \sqrt{\pi f \mu \sigma}$$

$$= \sqrt{\pi \times 10 \times 10^3 \times \mu_0 \times 10^{-3}}$$

$$= \sqrt{\pi \times 10 \times 4\pi \times 10^{-10} \times 10^3}$$

$$= \sqrt{40\pi^2 \times 10^{-7}}$$

$$\alpha = 6.28 \times 10^{-3} \text{ //}$$

$$\beta = 6.28 \times 10^{-3} \text{ //}$$

$$\gamma = \alpha + j\beta$$

$$= 6.28 \times 10^{-3} + j(6.28 \times 10^{-3})$$

$$= 8.88 \times 10^{-3} \angle 45^\circ$$

✓ Polar form
✓ ✓

$$\eta = \sqrt{\frac{\omega \mu}{\sigma}} \angle 45^\circ$$

$$= \sqrt{\frac{20\pi \times 10^3 \times 4\pi \times 10^{-7} \times 10^{-3}}{10^3}} \angle 45^\circ$$

$$= \sqrt{\quad}$$