

An alternating quantity can be represented using

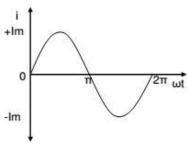
- i) Waveform
- ii) Equations
- iii) Phasor

A sinusoidal alternating quantity can be represented by a rotating line called a **Phasor.** A phasor is a line of definite length rotating in anticlockwise direction at a constant angular velocity

The waveform and equation representation of an alternating current is as shown. This sinusoidal quantity can also be represented using phasors.



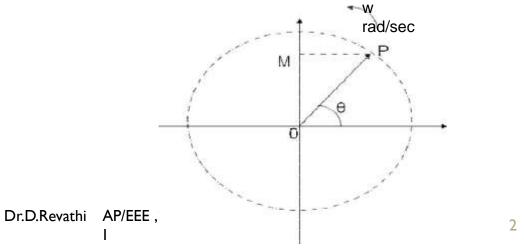




 $i = I_m \sin \omega t$

In phasor form the above wave is written as $\bar{I} = I_m \angle 0^0$

Draw a line OP of length equal to I_m . This line OP rotates in the anticlockwise direction with a uniform angular velocity ω rad/sec and follows the circular trajectory shown in figure. At any instant, the projection of OP on the y-axis is given by OM=OPsin $\theta = I_m sin\omega t$. Hence the line OP is the phasor representation of the sinusoidal current.



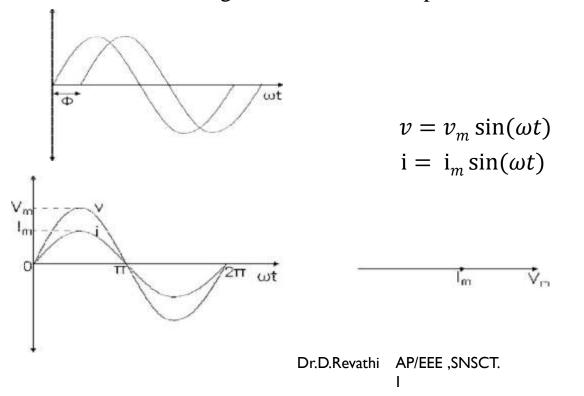




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In phase –

Two waveforms are said to be **in phase**, when the phase difference between them is zero. That is the zero points of both the waveforms are same. The waveform, phasor and equation representation of two sinusoidal quantities which are in phase is as shown. The figure shows that the voltage and current are in phase.

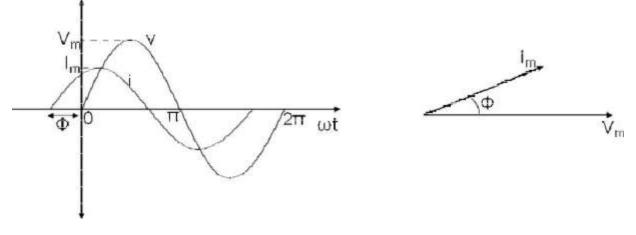






Leading

In the figure shown, the zero point of the current waveform is before the zero point of the voltage waveform. Hence the current is leading the voltage. The waveform, phasor and equation representation is as shown



$$v = v_m \sin(\omega t) = \overline{V}_m \angle 0^0$$

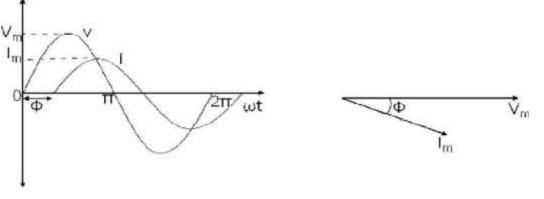
i = i_m sin(\omega t + \theta) = \overline{V}_m \angle \theta^0





lagging

In the figure shown, the zero point of the current waveform is after the zero point of the voltage waveform. Hence the current is lagging behind the voltage. The waveform, phasor and equation representation



$$v = v_m \sin(\omega t) = \stackrel{-}{=} V_m \angle 0^0$$

i = i_m sin(\omega t - \theta) = $\stackrel{-}{=} I_m \angle -\theta$



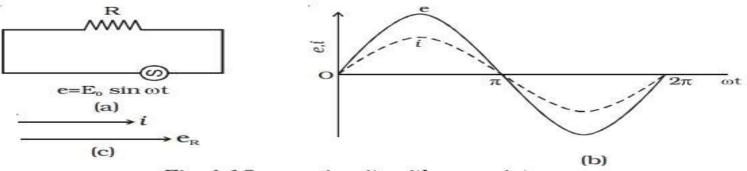
AC Circuit with resistor, inductor and capacitor

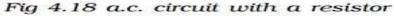


AC Circuit with resistor

Let an alternating source of emf be connected across a resistor of resistance R. The instantaneous value of the applied emf is

 $e = E_0 \sin \omega t$. . .(1)





If i is the current through the circuit at the instant t, the potential drop across R is, e = i R

Potential drop must be equal to the applied emf.

Hence, $iR = E_0 \sin \omega t$

$$i = \frac{E_o}{R} \sin \omega t$$
; $i = I_o \sin \omega t$...(2)
where $I_o = \frac{E_o}{R}$, is the peak value of a.c in the circuit. Equation

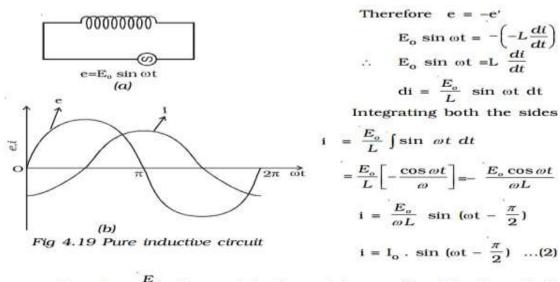
In a resistive circuit, the applied voltage and current are in phase with each other

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a negligible resistance and is wound on a laminated iron core. Due to an alternating emf that is applied to the inductive coil, a self induced emf is generated which opposes the applied voltage. (eg) Choke coil.

Let an alternating source of emf be applied to a pure inductor of inductance L. The inductor has

Inductive and Capacitive elements



AC Circuit with an inductor

where $I_0 = \frac{E_a}{\omega L}$. Here, ωL is the resistance offered by the coil. It is called inductive reactance. Its unit is ohm .

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From equations (1) and (2) it is clear that in an a.c. circuit containing a pure inductor the current i lags behind the voltage e phase angle of $\pi/2$.

Conversely the voltage across L leads the current by the phase angle of $\pi/2$. This fact is presented graphically in Fig 4.19b.

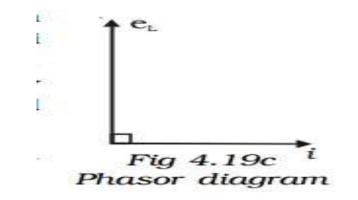


Fig 4.19c represents the phasor diagram of a.c. circuit containing only L. *Inductive reactance*

 $X_L = \omega L = 2\pi v L$, where v is the frequency of the a.c. supply For d.c. v = 0; $\therefore X_L = 0$

Thus a pure inductor offers zero resistance to d.c. But in an a.c. circuit the reactance of the coil increases with increase in freque

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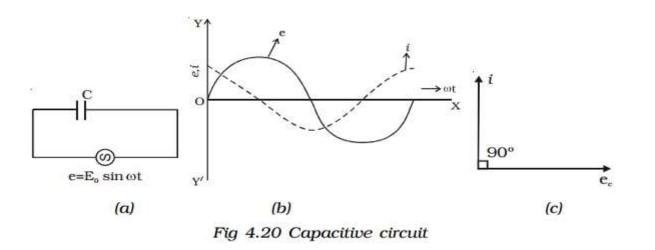


Pure Capacitive



AC Circuit with a capacitor

An alternating source of emf is connected across a capacitor of capacitance C (Fig 4.20a). It is charged first in one direction and then in the other direction.



Pure Capacitive circuit



The instantaneous value of the applied emf is given by

 $e = E_0 \sin \omega t$...(1)

At any instant the potential difference across the capacitor will be equal to the applied emf

 \therefore e = q/C, where q is the charge in the capacitor

But

$$i = \frac{dq}{dt} = \frac{d}{dt} (Ce)$$

$$i = \frac{d}{dt} (C E_{o} \sin \omega t) = \omega CE_{o} \cos \omega t$$

$$i = \frac{E_{o}}{(1/\omega C)} \sin\left(\omega t + \frac{\pi}{2}\right)$$

$$i = I_{o} \sin\left(\omega t + \frac{\pi}{2}\right) \dots (2)$$

where $I_0 = \frac{E_o}{(1/\omega C)}$ $\frac{1}{\omega C} = X_C$ is the resistance offered by the capacitor. It is called

capacitive reactance. Its unit is ohm .



Pure Capacitive circuit



From equations (1) and (2), it follows that in an a.c. circuit with a capacitor, the current leads the voltage by a phase angle of $\pi/2$. In other words the emf lags behind the current by a phase angle of $\pi/2$. This is represented graphically in Fig 4.20b.

Fig 4.20c represents the phasor diagram of a.c. circuit containing only C.

$$X_{\rm C} = \frac{1}{\omega C} = \frac{1}{2\pi \nu C}$$

where v is the frequency of the a.c. supply. In a d.c. circuit v = 0

$$\therefore$$
 $X_{\rm C} = \infty$

Thus a capacitor offers infinite resistance to d.c. For an a.c. the capacitive reactance varies inversely as the frequency of a.c. and also inversely as the capacitance of the capacitor.